

In this expression \hat{T} places each term of $S(\beta)$ in the correct place between $\psi(\tau)$, $\psi^+(0)$, and the multiplication rule for $S(\tau_2, \tau_1)$ has been used.

The perturbation theory is now developed as an expansion in powers of V_I . This requires to evaluate expectation values of higher-order products of field operators with equal numbers of ψ and of ψ^+ .

This is substantially facilitated by a factorization rule, Wick's theorem.

1.4.2 Wick's theorem

We define the many-particle Green's function for a non-interacting system as

$$G_n^0(\tau_1, \dots, \tau_n; \tau_1', \dots, \tau_n') = -\frac{1}{Z_G^0} \text{tr} \left\{ e^{-(H_0 - \mu N)\beta} \hat{T} \left\{ \psi(\tau_1) \dots \psi(\tau_n) \psi^+(\tau_1') \dots \psi^+(\tau_n') \right\} \right\}$$

where all operators are in the interaction picture:

$$\psi(\tau_i) = e^{(H_0 - \mu N)\tau_i} \psi e^{-(H_0 - \mu N)\tau_i} \quad \text{etc.}$$

The factorization of the tr of operator products is derived using the equation of motion for G :

The EOM for the thermal single-particle Green's fct. is

$$\left[-\frac{\partial}{\partial \tau} - (H_0 - \mu N)_{\tau} \right] G(\tau, \tau') = \delta(\tau - \tau')$$

← acting on the operator at time τ

The EOM for the non-interacting n -particle Green's function is then obtained as

$$\begin{aligned} & \left[-\frac{\partial}{\partial \tau_i} - (H_0 - \mu N)_{\tau_i} \right] G_n^0(\tau_1, \dots, \tau_n; \tau'_1, \dots, \tau'_n) \\ &= \sum_{j=1}^n (\mp 1)^{j-1+n-i} \delta(\tau_i - \tau'_j) G_{n-1}^0(\tau_1, \dots, \cancel{\tau_i}, \dots, \tau_n; \tau'_1, \dots, \cancel{\tau'_j}, \dots, \tau'_n) \end{aligned}$$

number of commutations needed to bring $\Psi^+(\tau'_j)$ and $\Psi(\tau_i)$ together.

$$\stackrel{\text{S.P. EOM}}{=} G_n^0(\tau_1, \dots, \tau_n; \tau'_1, \dots, \tau'_n) = \sum_{j=1}^n (\mp 1)^{j-1+n-i} G(\tau_i, \tau'_j) G_{n-1}^0(\tau_1, \dots, \cancel{\tau_i}, \dots, \tau_n; \tau'_1, \dots, \cancel{\tau'_j}, \dots, \tau'_n)$$

By induction it follows Wick's theorem:

The non-interacting n -particle Green's function (of operators with canonical commutation relations!) factorizes into a sum over products of free 1-particle Green's functions $G_1^0(\tau_i, \tau'_j)$. The sum runs over all possible pairings of operators $\Psi(\tau_i), \Psi^+(\tau'_j)$. For each pair there is a factor $(\mp 1)^P$ (fermions/bosons) where P is the number of elementary commutations required to bring the respective operators together.